

Solutions I

1 Convert the vector expression to suffices:

$$\begin{aligned} t_i &= -pn_i + \mu \left(2n_j \frac{\partial u_i}{\partial x_j} + \epsilon_{ijk} n_j \epsilon_{klm} \frac{\partial u_m}{\partial x_l} \right) \\ &= -pn_i + \mu \left(2n_j \frac{\partial u_i}{\partial x_j} + (\delta_{il} \delta_{jm} - \delta_{im} \delta_{jl}) n_j \frac{\partial u_m}{\partial x_l} \right) \\ &= -pn_i + \mu \left(2n_j \frac{\partial u_i}{\partial x_j} + n_j \frac{\partial u_j}{\partial x_i} - n_j \frac{\partial u_i}{\partial x_j} \right). \end{aligned}$$

2 The components of τ_{ij} are

$$\tau_{ij} = \begin{pmatrix} 0 & -2Uy/a^2 & -2Uz/b^2 \\ -2Uy/a^2 & 0 & 0 \\ -2Uz/b^2 & 0 & 0 \end{pmatrix}.$$

The normal to the plane $x = 0$ has components $(1, 0, 0)$, so the matrix-vector product gives the shear stress $(0, -2Uy/a^2, -2Uz/b^2)$. The magnitude of the stress is hence $f = 2U(y^2/a^4 + z^2/b^4)^{1/2}$. This is smallest on the axis where $y = z = 0$ and largest on the boundary, where we write $y = a \cos \theta$, $z = b \cos \theta$. Then

$$\frac{f^2}{4U^2} = \frac{\cos^2 \theta}{a^2} + \frac{\sin^2 \theta}{b^2} = \frac{1}{a^2} + \left(\frac{1}{b^2} - \frac{1}{a^2} \right) \sin^2 \theta = \frac{1}{b^2} + \left(\frac{1}{a^2} - \frac{1}{b^2} \right) \cos^2 \theta.$$

If $a < b$ this is maximum when $\theta = 0$ and $\theta = \pi$. If $a > b$ this is maximum when $\theta = \pi/2$ and $\theta = 3\pi/2$. So the shear stress magnitude is largest at the intersection of the semi-minor axis and the boundary. If $a = b$ (circle), the shear stress has the same value over the whole boundary.

3 Along any circle $\mathbf{n} = \hat{\mathbf{r}}$ so

$$\mathbf{u} \cdot \mathbf{n} = \nabla \phi \cdot \hat{\mathbf{r}} = \frac{\partial \phi}{\partial r} = U \cos \theta \left(1 - \frac{a^2}{r^2} \right).$$

This vanishes along the circle $r = a$. The function p is given along $r = a$ by

$$-\frac{1}{2}\rho \left(4U^2 \sin^2 \theta + \frac{\Gamma^2}{4\pi^2 a^2} - \frac{2\Gamma U}{\pi a} \sin \theta \right).$$

Along the circle we have $d\mathbf{S} = a(\cos \theta, \sin \theta) d\theta$. Hence

$$-\int_S p d\mathbf{S} = \frac{1}{2}\rho \int_0^{2\pi} \left(4U^2 \sin^2 \theta + \frac{\Gamma^2}{4\pi^2 a^2} - \frac{2\Gamma U}{\pi a} \sin \theta \right) (\cos \theta, \sin \theta) a d\theta.$$

This gives $-\rho\Gamma U j$. Along any circle $\mathbf{x} = \mathbf{n}$ so $\mathbf{x} \times \mathbf{n} = \mathbf{0}$. Therefore the last integral vanishes.

4 The equation for the streamlines is $dx/u = dy/v$, which leads to

$$\frac{1+t}{x} dx = \frac{2+t}{2y} dy \quad \text{i.e.} \quad |x|^{1+t} = A|y|^{1+t/2}.$$

Along the pathlines

$$\frac{dx}{dt} = \frac{x}{1+t}, \quad \frac{dy}{dt} = \frac{2y}{2+t}.$$

Using the initial condition, this gives in parametric form

$$x = 1 + t, \quad y = \left(\frac{2+t}{2} \right)^2$$

for $t \geq 0$. Solving for $y(x)$ gives

$$y = \left(\frac{1+x}{2} \right)^2$$

for $x \geq 1$. Write t_* for the time at which the particle passes through $(1, 1)$. Then at $t = 1$

$$x = \frac{2}{1+t_*}, \quad y = \left(\frac{3}{2+t_*} \right)^2$$

for $0 \leq t_* \leq 1$. Solving for $y(x)$ gives

$$y = \left(\frac{3x}{2+x} \right)^2$$

for $1 \leq x \leq 2$.