

① 4.10

From the example in the book,

$$d\vec{E} = \frac{\rho_e b}{4\pi\epsilon_0} d\phi \frac{-\hat{r}b + \hat{z}h}{(b^2 + h^2)^{3/2}} \bigg|_{h=0} = \frac{\rho_e b}{4\pi\epsilon_0} \frac{-\hat{r}}{b^3} d\phi$$

to find $\vec{E}$ we must integrate, but it is important to note that $\hat{r}$ is not a constant in the integration, as the direction changes as ϕ changes. We must use the relation $\hat{r} = \hat{x}\cos\phi + \hat{y}\sin\phi$

$$\vec{E} = -\frac{\rho_e b^2}{4\pi\epsilon_0} \cdot \frac{1}{b^3} \int_0^\pi (\hat{x}\cos\phi + \hat{y}\sin\phi) d\phi$$

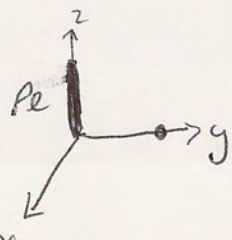
The $\hat{x}$ component vanishes, as expected from the symmetry of the problem.

$$\vec{E} = \frac{-\rho_e}{4\pi\epsilon_0 b} \hat{y} [-\cos\phi]_{\phi=0}^{\phi=\pi}$$

$$= -\frac{\rho_e}{2\pi\epsilon_0 b} \hat{y}$$

② 4.14

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{\rho_e dl'}{|\vec{r} - \vec{r}'|^3} (\vec{r} - \vec{r}')$$



where $\vec{r} = 10\text{cm} \hat{y}$ = observation point

and $\vec{r}' = z' \hat{z}$ = vector pointing to charge distribution

② 4.14

$dl' = dz' =$ differential length along line charge

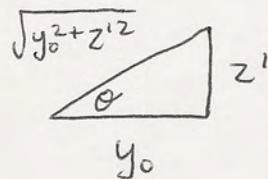
$$\text{so } R = |\vec{r} - \vec{r}'| = \sqrt{z'^2 + (.1\text{m})^2}$$

$$\vec{E} = \frac{\rho_L}{4\pi\epsilon_0} \int_{z'=0}^{z'=.05\text{m}} \frac{-z' \hat{z} + .1\text{m} \hat{y}}{[z'^2 + (.1\text{m})^2]^{3/2}} dz' = \underbrace{\frac{\rho_L}{4\pi\epsilon_0} \int_{z'=0}^{z'=.05\text{m}} \frac{-z' \hat{z} dz'}{(z'^2 + .1\text{m}^2)^{3/2}}}_{E_z} + \underbrace{\frac{\rho_L}{4\pi\epsilon_0} \int_{z'=0}^{z'=.05\text{m}} \frac{+.1\text{m} \hat{y} dz'}{(z'^2 + .1\text{m}^2)^{3/2}}}_{E_y}$$

These two integrals require a trigonometric substitution

let $y_0 = .1\text{m}$. From the triangle drawn,

$$\tan \theta = \frac{z'}{y_0} \Rightarrow z' = y_0 \tan \theta$$



$$dz' = y_0 \sec^2 \theta d\theta$$

$$\sin \theta = \frac{z'}{\sqrt{z'^2 + y_0^2}}, \quad \cos \theta = \frac{y_0}{\sqrt{z'^2 + y_0^2}}$$

Therefore

$$E_z = \frac{-\rho_L}{4\pi\epsilon_0} \int_{z'=0}^{z'=.05\text{m}} \sin \theta \cdot \frac{\cos^2 \theta}{y_0^2} \sec^2 \theta d\theta = \frac{-\rho_L}{4\pi\epsilon_0} \int \frac{\sin \theta d\theta}{y_0^2}$$

$$= + \frac{\rho_L}{4\pi\epsilon_0 y_0^2} \cos \theta = + \frac{\rho_L}{4\pi\epsilon_0 y_0^2} \frac{y_0}{\sqrt{z'^2 + y_0^2}} \bigg|_{z'=0\text{m}}^{z'=.05\text{m}} = -7.59 \times 10^5 \frac{\text{V}}{\text{m}}$$

$$E_y = \frac{+\rho_L}{4\pi\epsilon_0} \int_{z'=0}^{z'=.05\text{m}} y_0 \frac{\cos^3 \theta}{y_0^3} \sec^2 \theta d\theta = \frac{+\rho_L}{4\pi\epsilon_0 y_0^2} \int \cos \theta d\theta$$

$$= \frac{+\rho_L}{4\pi\epsilon_0 y_0^2} \sin \theta = \frac{\rho_L}{4\pi\epsilon_0 y_0^2} \frac{z'}{\sqrt{z'^2 + y_0^2}} \bigg|_{z'=0\text{m}}^{z'=.05\text{m}} = 3.22 \times 10^6 \frac{\text{V}}{\text{m}}$$

$$\vec{E}(0, .10\text{m}, 0) = -7.59 \times 10^5 \frac{\text{V}}{\text{m}} \hat{z} + 3.22 \times 10^6 \frac{\text{V}}{\text{m}} \hat{y}$$

3 4.24

Choose a sphere as a surface to integrate over

$$\oint \vec{E} \cdot d\vec{S} = \frac{1}{\epsilon_0} Q_{\text{enclosed}}$$

By symmetry, $\vec{E} = E_R \hat{R}$, where E_R depends only on R

$$d\vec{S} = R \sin\theta \, d\phi \, d\theta \, \hat{R}$$

$$\begin{aligned} \text{so } \oint \vec{E} \cdot d\vec{S} &= \int_0^\pi d\theta \int_0^{2\pi} d\phi \, E_R R \sin\theta \\ &= 4\pi R E_R \end{aligned}$$

Now, for $R < a$, $Q_{\text{enc}} = 0$, so

$$4\pi R E_R = 0$$

$$\boxed{E_R = 0}$$

For $a < R < b$, $Q_{\text{enc}} = Q_1$

$$4\pi R E_R = \frac{Q_1}{\epsilon_0}$$

$$E_R = \frac{Q_1}{4\pi\epsilon_0 R} \Rightarrow \vec{E} = \frac{Q_1}{4\pi\epsilon_0 R} \hat{R}$$

For $R > b$, $Q_{\text{enc}} = Q_1 + Q_2$

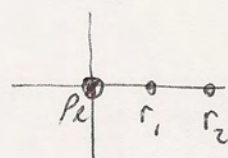
$$E_R = \frac{Q_1 + Q_2}{4\pi\epsilon_0 R} \Rightarrow \vec{E} = \frac{Q_1 + Q_2}{4\pi\epsilon_0 R} \hat{R}$$

4.33

The electric field from an infinite line charge is

$$\vec{E} = \frac{\rho_L}{2\pi\epsilon_0 r} \hat{r}$$

Because the fields depend only on r , the two points at radial distances r_1 and r_2 may be chosen as colinear with the line charge. The potential difference is



$$\Delta V = - \int_{r_1}^{r_2} \vec{E} \cdot d\vec{\ell} \quad d\vec{\ell} = \hat{r} dr$$

$$= - \frac{\rho_L}{2\pi\epsilon_0} \int_{r_1}^{r_2} \frac{1}{r} dr$$

$$= - \frac{\rho_L}{2\pi\epsilon_0} \ln \frac{r_2}{r_1}$$

⑤ 4.35

$$A + \theta = 0^\circ, R = 1m,$$

$$|\vec{E}| = \frac{qd}{4\pi\epsilon_0 (1m)^3} \cdot 2 = 4 \frac{mV}{m}$$

$$A + R = 2, \theta = 90^\circ$$

$$\vec{E} = \frac{qd}{4\pi\epsilon_0 (2m)^3} \hat{\theta}$$

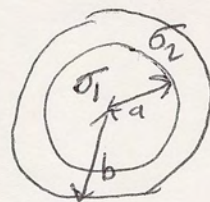
$$= \frac{4mV}{m} \cdot \frac{1}{16} \hat{\theta}$$

$$= .25 \frac{mV}{m} \hat{\theta}$$

$$= -.25 \frac{mV}{m} \hat{z}$$

⑥ 4.44

This can be modeled as two resistors in parallel:



$$R = R_1 // R_2 = \frac{l}{\sigma_1 A_1} // \frac{l}{\sigma_2 A_2}$$

$$A_1 = \pi a^2$$

$$A_2 = \pi(b^2 - a^2)$$

$$R = \frac{\frac{l^2}{\sigma_1 \sigma_2 A_1 A_2}}{\frac{l}{\sigma_1 A_1} + \frac{l}{\sigma_2 A_2}} = \frac{l}{\sigma_2 A_2 + \sigma_1 A_1}$$

$$R = \frac{l}{\pi [\sigma_1 a^2 + (b^2 - a^2) \sigma_2]}$$

⑦ 4.51 In free-space, $E_{t0} = E_{n0}$, so the angle is 45°

In medium 1, $E_{t1} = E_{t0}$, but $E_{n1} = \frac{\epsilon_0 E_{n0}}{3\epsilon_0} = \frac{1}{3} E_{n0}$

$$\text{So } \theta_1 = \tan^{-1}\left(\frac{E_{t1}}{E_{n1}}\right) = \tan^{-1}\left(\frac{E_{t0}}{\frac{1}{3}E_{n0}}\right) = \tan^{-1}(3) = 71.56^\circ$$

In medium 2, $E_{t2} = E_{t0}$, $E_{n2} = \frac{3\epsilon_0 E_{n1}}{5\epsilon_0}$

$$\theta_2 = \tan^{-1}\left(\frac{E_{t2}}{E_{n2}}\right) = \tan^{-1}\left(\frac{E_{t1}}{\frac{3}{5}E_{n1}}\right) = \tan^{-1}(5) = 78.69^\circ$$

In medium 3, following the same procedure,

$$\theta_3 = \tan^{-1}(7) = 81.87^\circ$$

8) 4.54

a) The electric field strength is $|\vec{E}| = \frac{V}{d} = \frac{10V}{.01m} = 10^3 \frac{V}{m}$
 $F = q|\vec{E}| = 1.6 \times 10^{-19} C \frac{V}{d} = 1.6 \times 10^{-16} N$

b) $a = \frac{F}{m} = \frac{1.6 \times 10^{-16} N}{9.1 \times 10^{-31} kg} = 1.76 \times 10^{14} \frac{m}{s^2}$

c) $d = \frac{1}{2} a t^2$
 $t = \sqrt{2d/a}$
 $= 1.07 \times 10^{-8} s$

Alternatively, use conservation of energy

$$qV = \frac{1}{2} m v_f^2 \Rightarrow v_f = \sqrt{\frac{2qV}{m}}$$
$$t = \frac{d}{\frac{1}{2}(v_f - 0)} = \frac{d}{\frac{1}{2} \sqrt{\frac{2qV}{m}}}$$
$$= 1.07 \times 10^{-8} s$$

9) 4.56

a) Boundary conditions require $E_{1t} = E_{2t}$, so
 $|\vec{E}_1| = |\vec{E}_2| = \frac{V}{d}$

b) $W_{e1} = \frac{1}{2} \int \epsilon_1 |\vec{E}|^2 dV = \frac{\epsilon_1}{2} A_1 d \left(\frac{V}{d}\right)^2 = \frac{1}{2} \frac{\epsilon_1 A_1}{d} V^2$
 $W_{e2} = \frac{1}{2} \int \epsilon_2 |\vec{E}|^2 dV = \frac{\epsilon_2}{2} A_2 d \left(\frac{V}{d}\right)^2 = \frac{1}{2} \frac{\epsilon_2 A_2}{d} V^2$

The energy stored in a capacitor is known to be
 $W_e = \frac{1}{2} C V^2$, so equating the expressions gives

$$C_1 = \frac{\epsilon_1 A_1}{d}, \quad C_2 = \frac{\epsilon_2 A_2}{d}$$

c) $W_{e1} + W_{e2} = \frac{1}{2} \left(\frac{\epsilon_1 A_1}{d} + \frac{\epsilon_2 A_2}{d} \right) V^2$

$$\text{So } C = \frac{\epsilon_1 A_1 + \epsilon_2 A_2}{d}$$
$$= C_1 + C_2$$

10 4.62

Locate image charges and observe how they move
when the real charge moves

