

Solutions III.

1 Consider a small cylindrical volume of fluid. Then from mass conservation ρAv is constant along the volume. Take the differential

$$Av d\rho + (\rho v) dA + (\rho A) dv = 0$$

and divide by ρAv . Newton II gives $F_1 - F_2 = -\rho A_1 v_1^2 + \rho A_2 v_2^2$, i.e. the change in F is given by

$$dF = -d(\rho Av^2) = -d(\dot{m}v) = -\dot{m} dv.$$

This is the force exerted on the fluid. The force exerted by the fluid is the opposite.

2 Mass conservation in the rectangular control volume gives

$$-\rho w U \delta + \dot{m} + \rho w \int_0^\delta u(y) dy = 0,$$

where \dot{m} is the mass flux out of the boundary layer. This gives

$$\dot{m} = \rho w U \delta - \rho w U \int_0^\delta (2y/\delta - y^2/\delta^2) dy = \rho w U \delta / 3.$$

So fluid is leaving the boundary layer. In numbers $\dot{m} = 1.24 \times 0.6 \times 30 \times 5 \times 10^{-3} / 3 = 0.0372 \text{ kg s}^{-1}$.

3 Newton II in the same control volume in the x -direction (along the boundary) gives (assuming no pressure gradient across the volume)

$$-\rho w U^2 \delta + U \dot{m} + \rho w \int_0^\delta u(y)^2 dy = F_x.$$

The tricky bit here is the second term: the momentum flux across the top of the boundary layer where the velocity is U . The force on the boundary is $-F_x$, i.e.

$$(2/3)\rho w U^2 \delta - \rho w U^2 \int_0^\delta (2y/\delta - y^2/\delta^2)^2 dy = (2/15)\rho w U^2 \delta.$$

This force is positive: there is a force in the direction of the flow on the plate. Its value is 0.4464 N. Note that is the actual force, not the force per unit area.

4 Steady frictionless flow. Force parallel to plate is from Newton II

$$F = U \cos \theta \dot{m},$$

since the flow is oriented at an angle θ to the plate. The horizontal components and vertical components of the flow are hence

$$F_x = U \cos^2 \theta \dot{m}, \quad F_y = U \sin \theta \cos \theta \dot{m}.$$

So $F_x = 2F_y$ implies $\cos \theta = 2 \sin \theta$, i.e. $\tan \theta = 1/2$, which gives $\theta = 26.6^\circ$.

5 Consider a control volume bounded by the radii r_1 and r_2 with depths and velocities h_1, h_2 and v_1, v_2 respectively. Conservation of mass gives

$$2\pi r_1 \rho h_1 v_1 = 2\pi r_2 \rho h_2 v_2.$$

For momentum, we need to take a wedge-shaped volume with angle $\Delta\theta$, otherwise the forces integrate to zero by symmetry. Neglect frictional effects so the only forces acting are pressure forces. Newton II in the radial direction then gives

$$\Delta\theta(r_1 \frac{1}{2} \rho g h_1^2 - r_2 \frac{1}{2} \rho g h_2^2) = \Delta\theta(-r_1 \rho h_1 v_1^2 + r_2 \rho h_2 v_2^2).$$

The $\Delta\theta$ terms cancel, but the r_i terms survive. We obtain the following two equations

$$r_1 h_1 v_1 = r_2 h_2 v_2, \quad \frac{1}{2} g r_1 h_1^2 - \frac{1}{2} g r_2 h_2^2 = -r_1 h_1 v_1^2 + r_2 h_2 v_2^2.$$

These two equations are unpleasant to solve. In the two-dimensional case, the r_i terms go away and progress can be made.