MAE105 Introduction to Mathematical Physics Spring Quarter 2015 http://web.eng.ucsd.edu/~sgls/MAE105_2015/

Homework 2

Due April 24, 2015.

1 Separation of Variables: mixed boundary conditions. Consider the heat equation,

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} \,,$$

with boundary conditions

$$k\frac{\partial u}{\partial x}(x=0,t)=0$$
, and $u(x=L,t)=0$,

and the linear initial condition

$$u(x,t=0)=a(x-L).$$

(a) What is the steady-state solution when $\partial u / \partial t = 0$?

(b) Use separation of variables to solve the initial value problem for u(x, t).

2 Separation of Variables: A pretty bad model for combustion. Consider the same heat equation as above, but with source *cu* (proportional to *u*), where *c* is a constant. This heat-dependent source term might (poorly) model a source of heat arising from, for example, a chemical reaction like combustion. The heat equation becomes

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2} + c u \,.$$

Assume insulating boundary conditions such that

$$\frac{\partial u}{\partial x}(x=0,t) = \frac{\partial u}{\partial x}(x=L,t) = 0.$$

(a) Use separation of variables to find the general solution as a series of cosines.

(b) What is the critical value of *c* for which u(x, t) can increase in time?

3 Time-dependent forcing: The sun heating the ocean. Let's try to model for how the sun heats the ocean surface. We use the boundary conditions

$$u(z=0,t) = u_0 e^{i\omega t}$$
 and $\frac{\partial u}{\partial z}(z \to -\infty, t) \to 0$,

where z = 0 is the ocean surface and as $z \to -\infty$ we are descending into the abyssal depths of the ocean. We propose to model the action of the sun as a source term in the heat equation by solving

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial z^2} + Q_0 \mathrm{e}^{z/\lambda} \mathrm{e}^{\mathrm{i}\omega t} \,.$$

- (a) Assume $\omega = 0$ and $\partial u / \partial t = 0$. Find the steady-state solution to the problem.
- (b) Now solve the problem with $\omega \neq 0$. Propose $u(z,t) = w(z)e^{i\omega t}$, then derive an equation for w(z). Solve this equation.
- (c) Write down the real part of your solution for w(z). This is the solution you would find if you replaced " $e^{i\omega t}$ " in the source term and boundary condition with $\cos(\omega t)$.

3 Laplace's equation in a square. Consider Laplace's equation in Cartesian coordinates in (x, y),

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$

with the boundary conditions

$$u(x = 0, y) = -1,$$
 $u(x, y = 0) = 0,$
 $u(x = L, y) = 1,$ $u(x, y = L) = 0.$

A sketch is given below.



Figure 1: "Sketch" for problem 4.

- (a) Use the principle of superposition and separation of variables to find u(x, y) which satisfies the governing equation and all boundary conditions.
- (b) What is the solution for u(x, y) when the boundary conditions at y = 0 and y = L are both changed to

$$\frac{\partial u}{\partial y} = 0?$$

Finding the solution should not require more than a line or two of calculation. *Hint: will the solution depend on y?*

4 Laplace's equation outside a disk. Consider Laplace's equation outside the disk with radius *a*. The domain thus extends from r = a to ∞ . Laplace's equation in polar coordinates is

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial u}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 u}{\partial \theta^2} = 0.$$

The boundary condition at r = a is

$$\frac{\partial u}{\partial r}(r=a,\theta)=1+2\sin\theta\,.$$

The boundary condition as $r \to +\infty$ is

$$\nabla u(r \to \infty, \theta) \to 0$$

A sketch is below.



Figure 2: "Sketch" for problem 5.

(a) Solve for $u(r, \theta)$. Hint: your solution will contain an undeterminable constant. This is because there are two solutions with no θ -dependence.